

#### The simulation of stratospheric discharges sustained by the secondary electrons from cosmic rays

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#### **Motivation**

How to describe the initiation and the development of upper-atmospheric lightning (transient luminous events, TLE)?

- Typical electric field values (e.g. in stratosphere) are much less than breakdown values;
- In contrast to tropospheric lightning, TLE have a very wide range of visual manifestations (blue jets, blue starters, sprites, elfs, gigantic jets, etc.);



E.M. Wescott et al. // J. Geo. Res.: Space Phys. 106 (A10), 2001, pp. 21549-21554

#### Runaway breakdown\* (RB) = TLE origin?

In atmospheric conditions runaway breakdown takes place at electric fields an order of magnitude weaker than those needed for normal breakdown in air, when following conditions are satisfied:

1) The value electric E field strength exceeds so-called "critical" value  $E_c$ 

$$E \ge E_c \equiv \frac{4\pi e^3 Z N_n}{mc^2} a, \qquad E_c \approx 2.2 \frac{kV}{cm} \cdot \frac{P}{P_n} << E_{br} \approx 23 \frac{kV}{cm}$$

2) The spatial scale L of the electric field exceeds the length of the avalanche growth  $I_a$ 

$$L ? \quad l_a \equiv \frac{m^2 c^4}{2\pi N_n Z e^4} \frac{E_c}{E}, \qquad l_a \approx 50 \ m \cdot \frac{E_c}{E} \cdot \frac{P_0}{P}$$

3) Fast seed electrons with energies above mare present

$$\varepsilon \ge \varepsilon_c \equiv \frac{mc^2}{2} \cdot \frac{E_c}{E}, \qquad \varepsilon_c \approx 2.5 MeV \cdot \frac{E_c}{E}$$

<sup>3</sup> \*A.V. Gurevich, and K.P. Zybin // UFN, 44(11), pp.1119–1140 (2001)



#### **Upper-atmospheric conditions for RB**

1) "Critical" field value falls exponentially with the attitude  $E_c \sim exp(-h)$  and the air conductivity above 20-30 km is high, so there is almost **no electric field** at 20-50 km attitudes.

However, the **positive lightning often violates** the electroneutrality condition providing the electric fields above critical value  $E > E_c$ 

2) The characteristic sizes of clouds and their parts *L* are **always** much greater than the avalanche scales  $I_a$ ;

3) The flux of **secondary cosmic ray seed electrons** is very large. At 4-8 km its attitude average flux density  $\underline{\square} \sim 10^3 \text{ m}^{-2} \text{ s}^{-1}$  Typical mean energies are equal to 0.1-10 MeV.



#### Aims

To perform the direct numerical simulation of large-scale gas discharge (here "blue jet") in terms of physical kinetics. Namely, we need to compute the EDF at the discharge initiation and early development stages;





#### **Kinetic model outline\***

- Electrons. Relativistic Boltzmann equation for EDF  $\gamma(p) \left( \frac{\partial f(\stackrel{r}{x}, \stackrel{r}{p}, t)}{\partial t} + \frac{\stackrel{r}{p}}{m\gamma(p)} \frac{\partial f(\stackrel{r}{x}, \stackrel{r}{p}, t)}{\partial x} - q \left\{ \stackrel{r}{E}(\stackrel{r}{x}, t) + \frac{\stackrel{r}{p} \times \stackrel{r}{B}}{m\gamma(p)} \right\} \frac{\partial f(\stackrel{r}{x}, \stackrel{r}{p}, t)}{\partial p} = Q_c \left[ (\stackrel{r}{x}, \stackrel{r}{p}, t) \right]$
- Ions. Continuity equations with drift-diffusion approximation (if needed)

$$\frac{\partial n_{+}}{\partial t} + \nabla \Gamma_{+} = \int_{-\infty}^{\infty} Q_{c} \left[ f\left( \begin{matrix} \mathbf{r} & \mathbf{r} \\ x, p, t \end{matrix} \right) \right] dp,$$

• Electromagnetic field. Maxwell's equations

$$\nabla \vec{E} = \frac{q}{\varepsilon_0} [n_+ - n_-], \quad \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}, \quad \nabla \times \vec{B} = \vec{j} + \varepsilon_0 \frac{\partial \vec{E}}{\partial t}, \quad \nabla \vec{B} = 0$$

where

$$n_{-}(\overset{\mathbf{r}}{x},t) = q \int_{-\infty}^{\infty} f(\overset{\mathbf{r}}{x},\overset{\mathbf{r}}{p},t) d\overset{\mathbf{r}}{p}, \quad \overset{\mathbf{r}}{j_{-}}(\overset{\mathbf{r}}{x},t) = \frac{q}{m} \int_{-\infty}^{\infty} f(\overset{\mathbf{r}}{x},\overset{\mathbf{r}}{p},t) \frac{\overset{\mathbf{r}}{p}}{\gamma(p)} d\overset{\mathbf{r}}{p}$$

\*A. V. Kozyrev, V. Y. Kozhevnikov, et al. // EPL, v. 114, no. 4 (2016), p. 45001
 A. V. Kozyrev, V. Y. Kozhevnikov, et al. // EPL, v. 112, no. 1 (2015), p. 15001

# **Model parameters**



#### 1D spherically symmetric

Pressure ~ 0.1 atm, temperature ~ 216 K

Interelectrode distance ~ 100 meters, cloud inhomogeneity size ~ 1 meter

Applied voltage  $U_{\text{max}} = 5 \text{ MV} (U_{\text{br}} = 50 \text{ MV})$ 

Initial electron mean energy  $\implies 1$  MeV, with the density n0 ~ 10<sup>9</sup> m<sup>-3</sup> entirely localized near cloud



#### Elastic & Inelastic cross-sections (for air\*)



Nelder approximation  $\sigma(\varepsilon) = \frac{\varepsilon + a}{b_0 + b_1(\varepsilon + a) + b_2(\varepsilon + a)^2}$ 

Elastic high-energy limit (Bethe extrapolation)

$$\sigma(\varepsilon); \ \frac{2Z}{\pi} \left(\frac{q^2}{4\varepsilon_0\varepsilon}\right)^2 \left(\frac{\varepsilon}{I_0}-1\right)$$

Inelastic high-energy limit (Bethe extrapolation)

$$\sigma(\varepsilon); \ \pi\left(rac{Zq^2}{4\pi\varepsilon_0\varepsilon}
ight)^2 \log\left(rac{arepsilon}{I_0}
ight)$$

<sup>8</sup> Y. Itikawa // J. Phys. Chem. Ref. Data, v. 35, no 1 (2006)

# **Numerical method**

- Direct numerical solution of Boltzmann equation by using semi-Lagrangian method with second order splitting in time (Strang splitting);
- ✓ Cubic spline (Cheng-Knorr) & Linear interpolation schemes;
- ✓ Quasi-uniform "semi-infinity" 500x10001 phase-space computational grid;
- ✓ Average time step is 1-3 picoseconds.

G. Strang // SIAM J. Numer. Anal., v. 5, no. 3 (1968), pp. 506-517
 <sup>9</sup> C.Z. Cheng, G. Knorr // J. Comp. Phys., v. 22, no. 3 (1976), pp. 330-351
 M. Lesur, Y. Idomura, and X. Garbet // Phys. Plasmas, 16, 092305 (2009)



# EDF (log scale), plasma density & electric field







*U*<sub>max</sub> = 5 MV, **1** = 1 MeV @ 0.1 atm

#### Discharge in a "subcritical" electric field

EDF at 0 ns (in logarithmic scale)





*U*<sub>max</sub> = 2 MV, = 1 MeV @ 0.1 atm





Within the framework of the kinetic theory we have obtained some agreement with asymptotic theory of runaway breakdown influence on the upper-atmospheric discharge formation and initiation. Model provides the energy spectrum of runaway electrons as well as the detailed timedependent discharge EDF structure.

Future work to be done:

- Investigate the influence of the initial EDF on the discharge initiation and its formation;
- Simulate the "electrodeless" TLE forms, e.g. "elfes";
- 2D & 3D hybrid models...



#### Instead of a conclusion: Apokampic discharge\* TLE origin &Runaway breakdown?

**Apokampic discharge** (from Greek  $a\pi \dot{o}$  - "off" and  $\kappa a\mu \pi \eta$  - "bend") - the stable plasma jet developing perpendicular to the bending point of the pulsed arc discharge channel between two electrodes.

Following the high-speed recording, it was discovered that apokamp channel consists of a set of plasma bullets moving with supersonic velocity of about 100–220 km/s.



#### Blue jet (18 km altitude with the tip at 35 km)

Apokamp

<sup>13</sup> \*E. A. Sosnin et al. // JETP Letters, vol. 103, no. 12, pp. 761–764, (2016) E. A. Sosnin et al. // EPJ D, vol. 71, no. 2, (2017)



#### Instead of a conclusion: Apokampic discharge





 1 – floating potential electrode; 2 – highvoltage electrode; 3 – pulsed-periodic arc discharge channel; 4 – bright process;

5 – apokamp channel.

20 420 300 150 30 pressure in Torr

<sup>14</sup> <sup>\*</sup>E. A. Sosnin et al. // JETP Letters, vol. 103, no. 12, pp. 761–764, (2016) E. A. Sosnin et al. // EPJ D, vol. 71, no. 2, (2017)



# Thank you!

