## Propagation of Cosmic Rays in Magnetized Protoplanetary Disks

### Soonyoung Roh

UNIST (Ulsan National Institute of Science and Technology), Korea

in Collaboration with Yuri I. Fujii and Marco Padovani

2-4 May 2018 Arcetri, Florence, Italy



Planets form in disks around young starts. Credit: S. Andrews (Harvard-Smithsonian CfA), ALMA (ESO/NAOJ/NRAO)

### ALMA's Best Image of a Protoplanetary disk



#### Egde-on PPD Orion Nebula



New images from the ALMA reveal neverbefore-seen details in the planet-forming disk around Sun-like star.

### Astrophysical Disk

Circular, coplanar and thin, Keplerian

Shearing, dissipative systems

- ✦ Shearing and differential rotation leading to viscous forces which cause radial motions.
- Dissipation: angular momentum transport outward, matter moves inward, energy liberated.
  => allow the accretion of gas on to the central star
- ★ Angular momentum transport is the key to the evolution of protoplanetary disks.

## The Magnetorotational instability (MRI)

Magnetorotational instability (MRI) is a promising mechanism that drives Magnetohydrodynamic (MHD) turbulence in accretion disk.

MHD turbulence is triggered and sustained in the disk by the MRI.



Balbus & Hawley 1991; Balbus & Hawley (1991)

### The Magnetorotational instability (MRI)

The MRI requires good coupling between the disk gas and the magnetic field.

Sufficient ionization degree is needed.



#### Umebayashi&Nakano 1981

### Previous work

In many studies on ionization in PPDs, ionization rate by cosmic rays,  $\zeta_{\rm C}$ , is estimated as

$$\zeta_{\rm C} = \frac{\zeta_{\rm CR}}{2} \left\{ \exp\left[-\frac{\chi(r,z)}{\chi_{\rm CR}}\right] + \exp\left[-\frac{\Sigma(r) - \chi(r,z)}{\chi_{\rm CR}}\right] \right\}$$

### Umebayashi&Nakano 1981

### Previous work

In many studies on ionization in PPDs, ionization rate by cosmic rays,  $\zeta_{\rm C}$ , is estimated as

 $\zeta_{\rm CR} \sim 10^{-17} \, {\rm s}^{-1} : \text{ ionization rate by cosmic rays in the interstellar space}$ (Black et al. 1990)  $\zeta_{\rm C} = \frac{\zeta_{\rm CR}}{2} \left\{ \exp\left[-\frac{\chi(r,z)}{\chi_{\rm CR}}\right] + \exp\left[-\frac{\Sigma(r) - \chi(r,z)}{\chi_{\rm CR}}\right] \right\}$ 

 $X_{CR} \sim 96 \text{ g/cm}^2$ : attenuation length of cosmic rays in the molecular cloud (Umebayashi&Nakano 1981)

### Umebayashi&Nakano 1981

### Previous work

In many studies on ionization in PPDs, ionization rate by cosmic rays,  $\zeta_{\rm C}$ , is estimated as

 $\zeta_{\rm CR} \sim 10^{-17} \, {\rm s}^{-1}$ : ionization rate by cosmic rays in the interstellar space (Black et al. 1990)  $\zeta_{\rm C} = \frac{\zeta_{\rm CR}}{2} \left\{ \exp\left[-\frac{\chi(r,z)}{\chi_{\rm CR}}\right] + \exp\left[-\frac{\Sigma(r) - \chi(r,z)}{\chi_{\rm CR}}\right] \right\}$  $X_{CR} \sim 96 \text{ g/cm}^2$ : attenuation length of cosmic rays (Umebayashi&Nakano 1981)  $\chi(r,z) = \int_{z}^{\infty} \rho(r,z) dz \quad [g \text{ cm}^{-2}]$   $\downarrow : \text{ surface density } r : \text{ radius}$ z: vertical direction *!* : density mid-plane

In many studies on ionization in PPDs, ionization rate by cosmic rays,  $\zeta_{\rm C}$ , is estimated as

 $\zeta_{\rm CR} \sim 10^{-17} \, {\rm s}^{-1}$ : ionization rate by cosmic rays in the interstellar space (Black et al. 1990)  $\zeta_{\rm C} = \frac{\zeta_{\rm CR}}{2} \left\{ \exp\left[-\frac{\chi(r,z)}{\chi_{\rm CR}}\right] + \exp\left[-\frac{\Sigma(r) - \chi(r,z)}{\chi_{\rm CR}}\right] \right\}$  $X_{CR} \sim 96 \text{ g/cm}^2$ : attenuation length of cosmic rays (Umebayashi&Nakano 1981)  $!_{+} = "_{-z}^{!} "(r, z) dl [g cm^{-2}]$ ! : surface density r : radius ! : density z: vertical direction I: particle travel distance

### Propagation of CRs

Solve relativistic equation of motion (4th-order Runge-Kutta Method) including energy loss equation

$$\mathbf{P} = \gamma m_i \mathbf{v} \qquad \frac{d\mathbf{P}}{dt} = q_i (\mathbf{E} + \frac{\mathbf{v} \times \mathbf{B}}{c}) + \text{energy loss}$$
  
energy loss eq.  
$$L_k(E_k) = -\frac{1}{n(\mathbf{H}_2)} \left( \frac{dE_k}{d\ell} \right) = -\frac{dE_k}{dN(\mathbf{H}_2)}$$
$$N(H_2) = \mathbf{n}(\mathbf{H}_2) \mathbf{d}$$
  
dl (cm]: path length  $E_k$  (eV): CR kinetic energy  
 $\mathbf{n}(\mathbf{H}_2)$  (cm<sup>-3</sup>): number density of the molecular  
hydrogen in the medium  
 $N(\mathbf{H}_2)$  (cm<sup>-2</sup>): column number density of molecular  
hydrogen along a particle trajectory  
 $\mathbf{M}$ . Padovani+2009



# Simulation setup

\*Test particle calculations of CRs using a snapshot of a MHD simulation of the MRI-active protoplanetary disk as the background.

\*Used Particle energy ranges (0.5 million particles)  $KE_0 = 10^6 eV, 10^7 eV, 10^8 eV, 10^9 eV, 10^{10} eV,$ and  $10^{11} eV$ 

\*Ideal-MHD simulation in a local box located at 1AU of a MMSN type of the disk T~ 280K, plasma beta =  $10^5$  @mid-pnace

\* Box size (Lx, Ly, Lz)=(±1H, ±1H, ±4H)

\* Initial position and velocity distributions are randomly generated.

mid plane

### Particle distribution







### Cosmic Ray Flux

The following definition of the differential cosmic ray flux J is used:

where dN is the number of particles in the energy range dE, passing through an area and solid angle elements dA and dQ in a time interval dt.(A. Dmitry et al. 2016)





# Summary and Conclusion

- Ionization rate plays an important role in setting chemical and physical fate of disks and ultimately that of planets form there.
- CRs are believed to be a primary ionization source in deep inside of a disk (e.q., Glassgold et al. 1997; Walsh et al. 2012) due to their deep penetration.
- Actual path of cosmic ray propagation is complicated due to gyro motion along turbulent magnetic field.
- -In order to quantitatively analyze the propagation of CRs considering energy loss processes in magnetized PPDs, we have performed test particle calculations of CRs .
- -Consider the local interstellar spectrum of CRs and calculated their trajectories in the disk.

-Ionization Rate @ each height: r elation to the ionization rate within the disk.