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Understanding Primordial Star Formation: Francesco's Contribution



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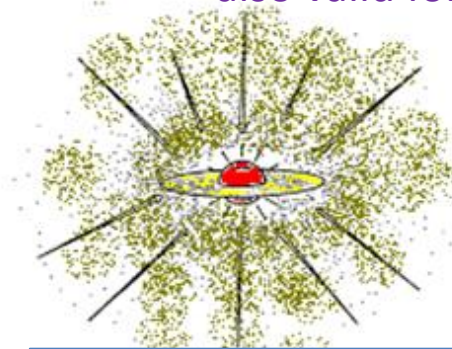
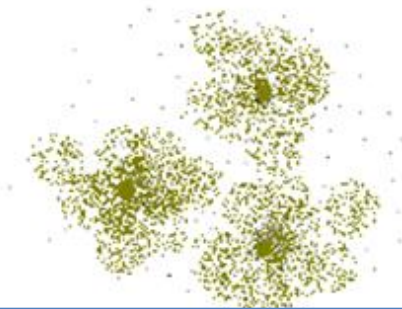
steps toward star formation

PSS/SPS papers (1983-86)

established in '80s

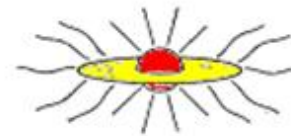
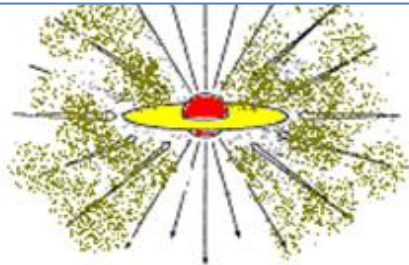
also valid for first star formation

Shu, Adams & Lizano (1987)



1) "Primordial star formation:
the role of molecular hydrogen"
Palla, Salpeter & Stahler (1983)

2) "Primordial stellar evolution:
the protostar phase"
Stahler, Palla & Salpeter (1986)

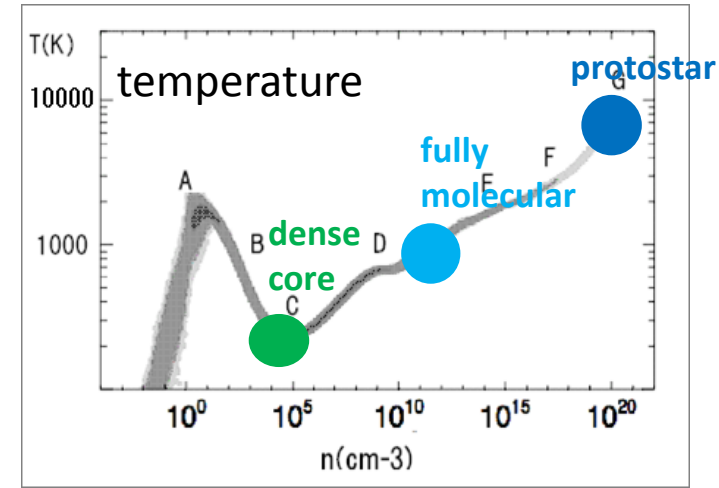
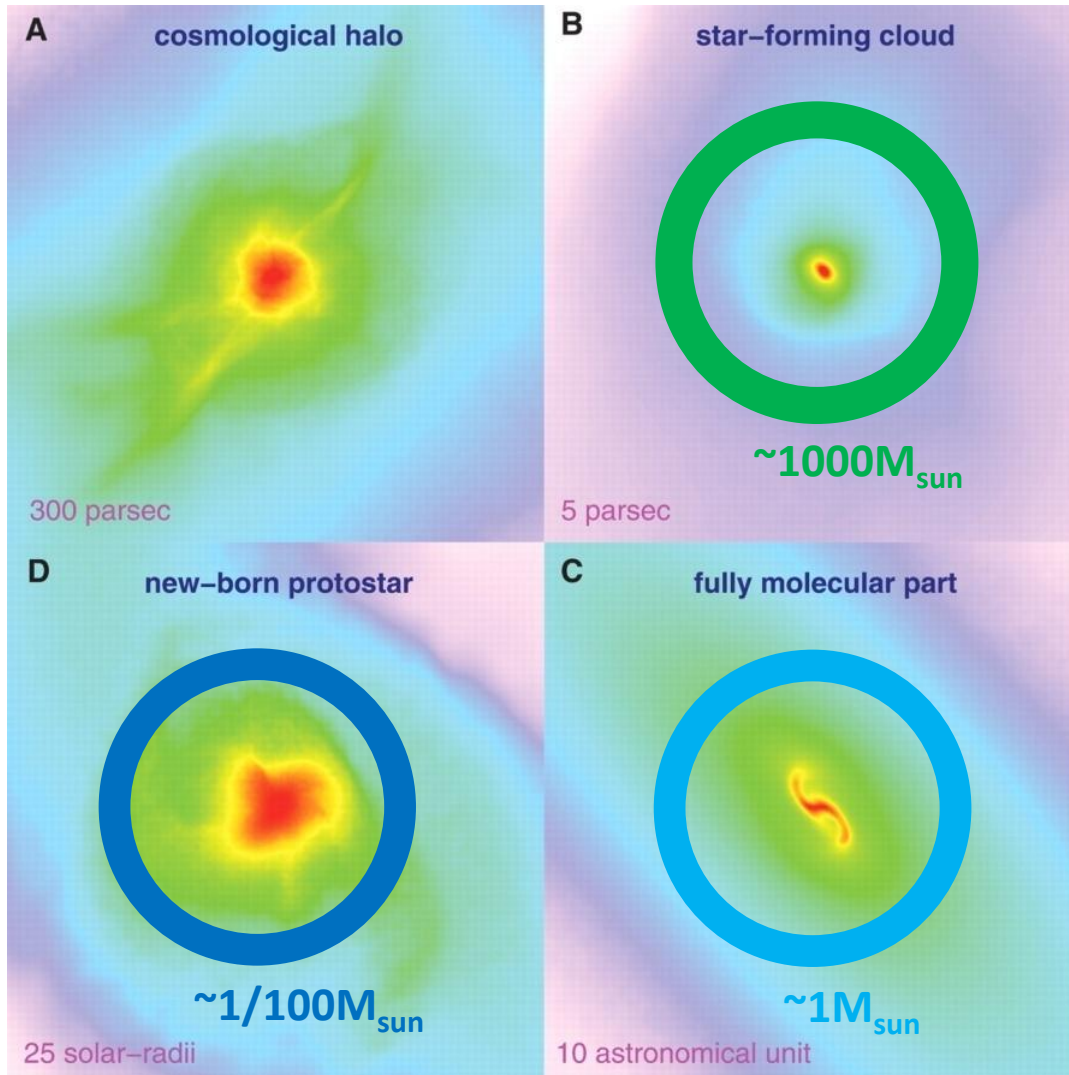


C) accretion stops by some
mechanism, e.g., by stellar
feedback

3) "Primordial stellar evolution:
the pre-main-sequence phase"
Stahler, Palla & Salpeter (1986)

Prestellar collapse of the first stars: current picture

Yoshida, KO, Hernquist 2008



Important physical scales:

- dense core ($\sim 1000 M_{\text{sun}}$) forms at $\sim 10^4 \text{cm}^{-3}$
- central $\sim 1 M_{\text{sun}}$ becomes fully H_2 at 10^{11}cm^{-3}
- small $\sim 10^{-2} M_{\text{sun}}$ hydrostatic protostar forms at $\sim 10^{21} \text{cm}^{-3}$

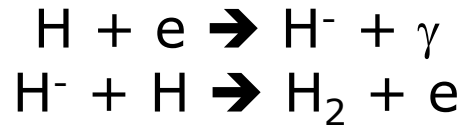
Prestellar collapse is controlled by H_2 cooling

H₂ formation in primordial gas

At low densities (<10⁸cm⁻³)

H⁻ channel : e catalyzed

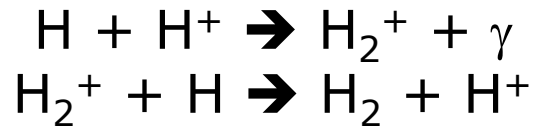
(Peebles & Dicke 1968; Hirasawa+1969)



dominant in low densities

H₂⁺ channel : H⁺ catalyzed

(Saslaw & Zipoy 1967)

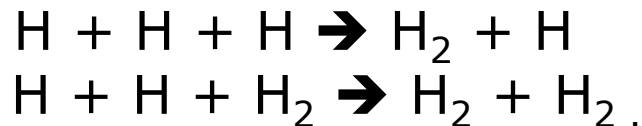


subdominant

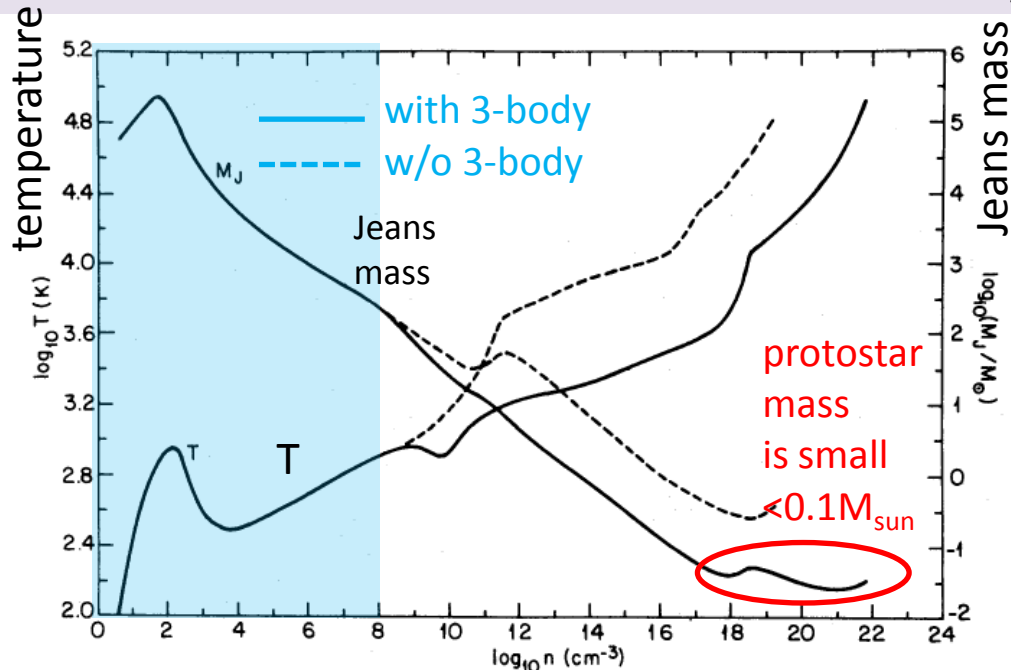
(can be important in strong radiation fields)

At higher densities (>10⁸cm⁻³)

3-body reactions (← Palla, Salpeter & Stahler 1983)



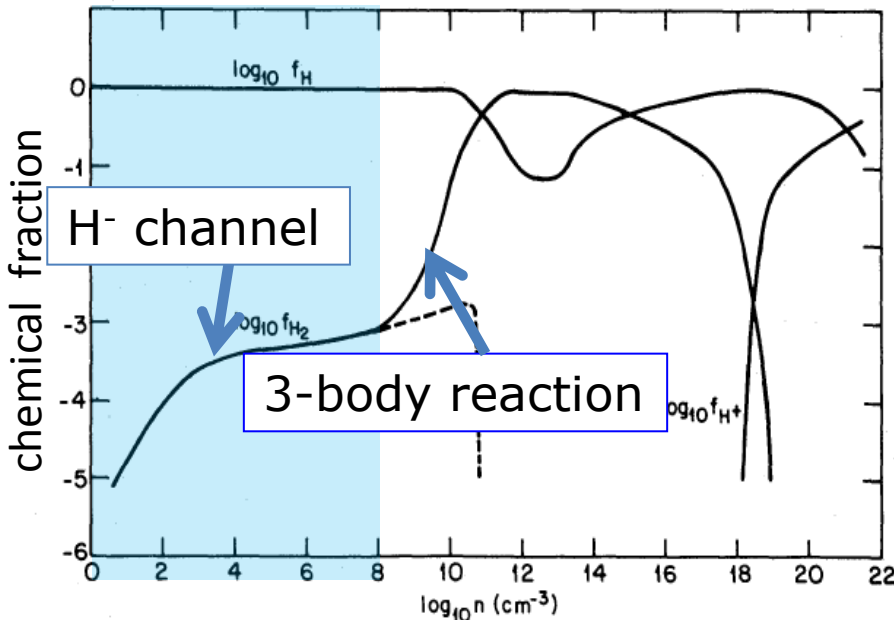
Effect of 3-body reaction



Palla, Salpeter & Stahler (1983)

In low densities, H_2 fraction reaches at most 10^{-3} by H^- channel.

Effect of 3-body reaction is dramatic at $> \sim 10^8 \text{cm}^{-3}$

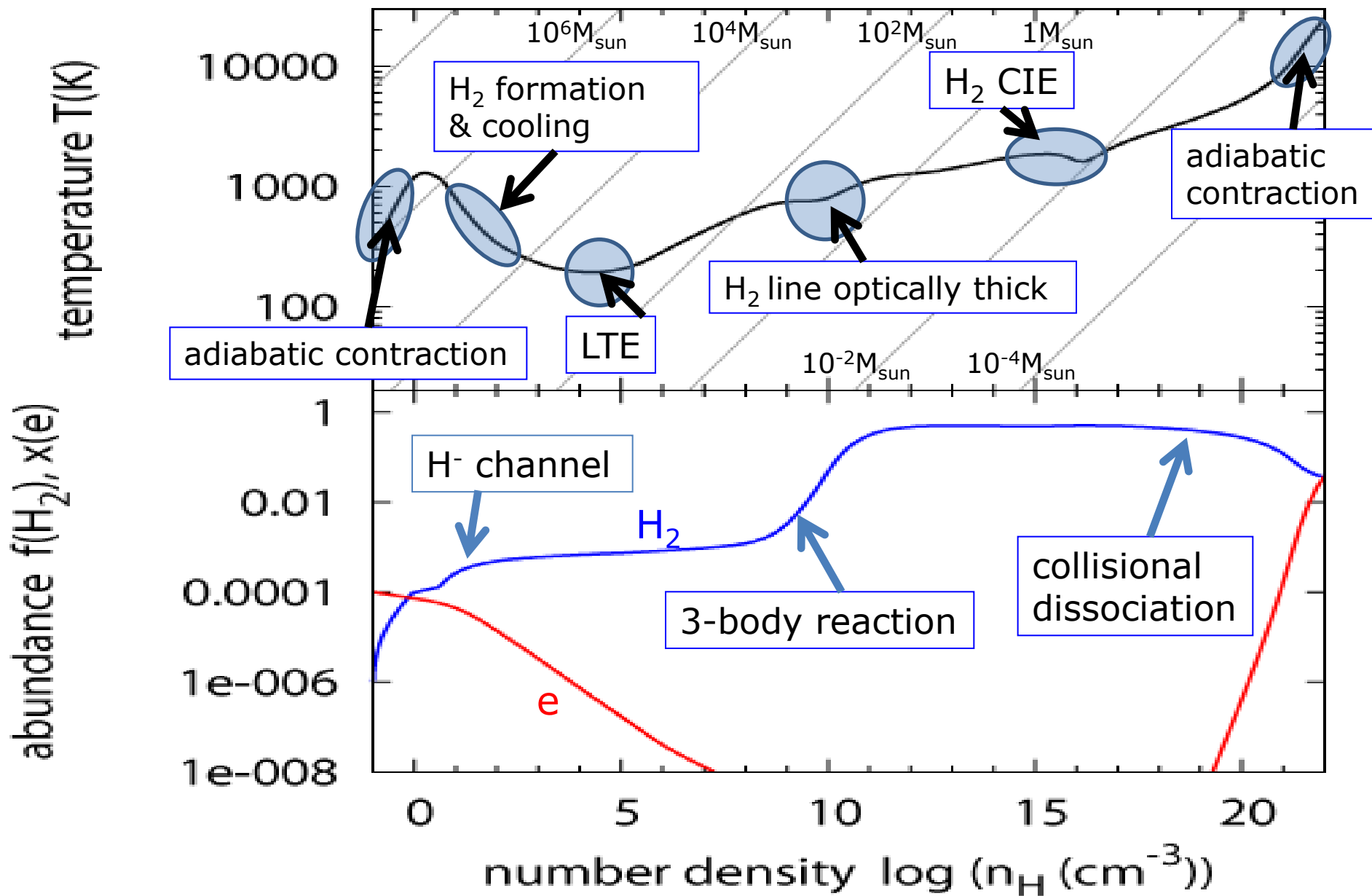


✓ all the hydrogen is converted to the molecular form.

✓ subsequent temperature evolution is largely altered.

✓ Jeans mass falls below $0.1 M_{\text{sun}}$. As a result, a small protostar will be formed.

Thermal evolution of primordial gas: current picture



H₂ CIE was already discussed

Palla, Salpeter & Stahler (1983)

APPENDIX

OMITTED REACTIONS AND RADIATIVE PROCESSES

Finally, we have considered the pressure-induced absorption of continuum photons in colliding H₂ molecules. Patch (1971) has computed the monochromatic absorption coefficient for this process, taking into account the excitation of rotational and vibrational levels. In a 50 M_{\odot} cloud ($n = 10^{16} \text{ cm}^{-3}$, $f_{\text{H}_2} = 1/3$, $T = 3000 \text{ K}$), photons near the blackbody peak see an optical depth of order unity from this process, although the optical depth is much lower for both lower densities (because of the quadratic density dependence of the opacity) and higher densities (because of the disappearance of the molecules). Using detailed balance, we find the emission rate from this process to be about 0.1 times the rate of compressional work on the gas at $n = 10^{16} \text{ cm}^{-3}$. It thus appears that this process, although it can become significant both for absorption and emission immediately prior to H₂ destruction, will not have a major effect on the thermal evolution.

In reality,

once H₂ CIE cooling becomes important,

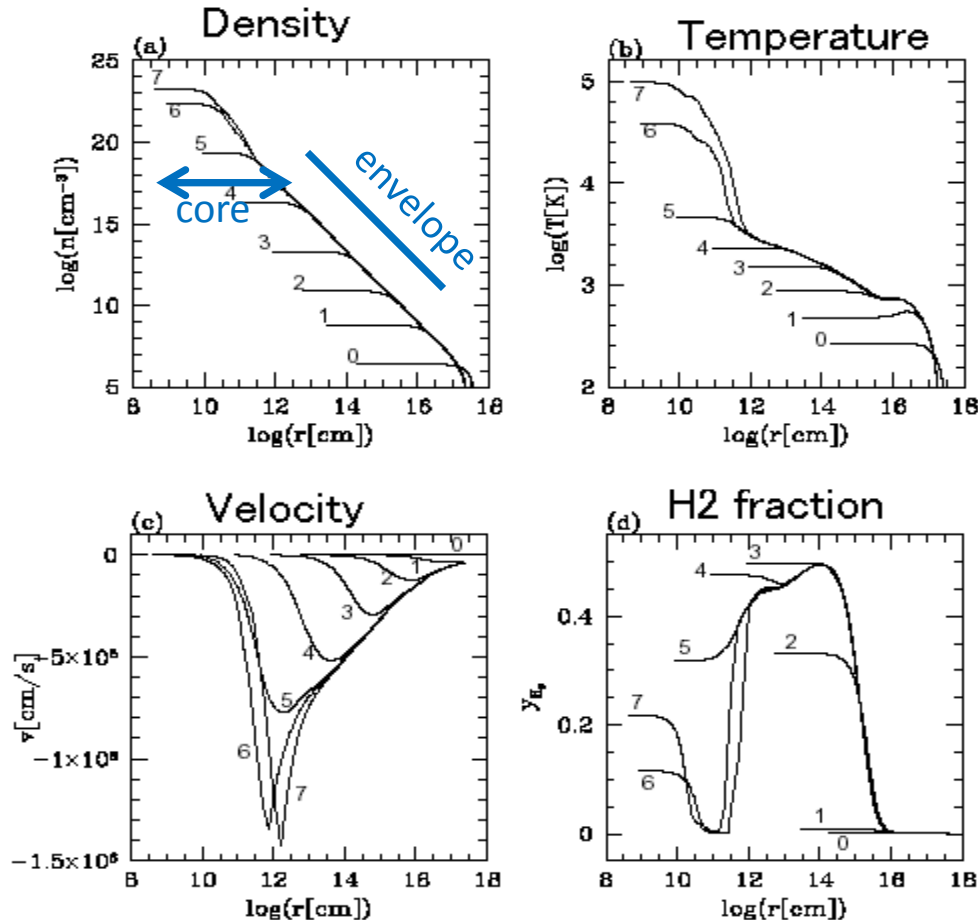
H₂ dissociation is delayed for other 2-3 orders of magnitude in density,

which results in further 1-2 orders of magnitude

reduction in formed protostellar mass.

Hydrodynamical evolution

1D radiation hydro calculation
KO & Nishi (1998)



core-envelope structure
develops

- envelope contains large mass but does not contribute so much to the optical depth
- effective mass of the cloud decreases during the collapse

It makes the cloud to cool more efficiently than the homogeneous contraction.

MOLECULAR HYDROGEN IN THE EARLY UNIVERSE

Firenze, 6-4 December 1997

Edited by E. Corbelli, D. Galli and F. Palla

Villa Agape



Francesco: "When and where did the Dark Age end?"

"We know the answer.

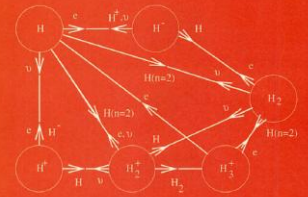
Here in Florence in 14th century!"

MEMORIE

DELLA SOCIETÀ ASTRONOMICA ITALIANA

JOURNAL OF THE ITALIAN ASTRONOMICAL SOCIETY

Vol. 69 - N. 2 - 1998



MOLECULAR HYDROGEN IN THE EARLY UNIVERSE

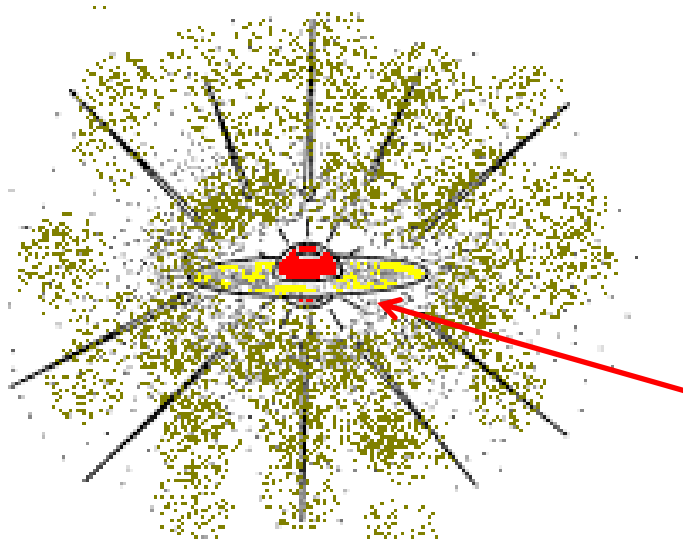
Firenze, 6-4 December 1997

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Protostellar evolution

SSP ('86)

accreting envelope
 $\sim 1000M_{\text{sun}}$



protostar
 $\sim 0.01M_{\text{sun}}$

mass accretion rate

$$\begin{aligned}\dot{M} &\cong M_J / t_{ff} \cong (c_s t_{ff})^3 \rho / t_{ff} \\ &\cong c_s^3 / G \propto T^{3/2}\end{aligned}$$

Pop III (1000K) $\sim 10^{-3}M_{\text{sun}}/\text{yr}$

Pop I (10K) $\sim 10^{-6}M_{\text{sun}}/\text{yr}$,

accretion rate much higher in Pop III case

Protostellar evolution by SPS('86)

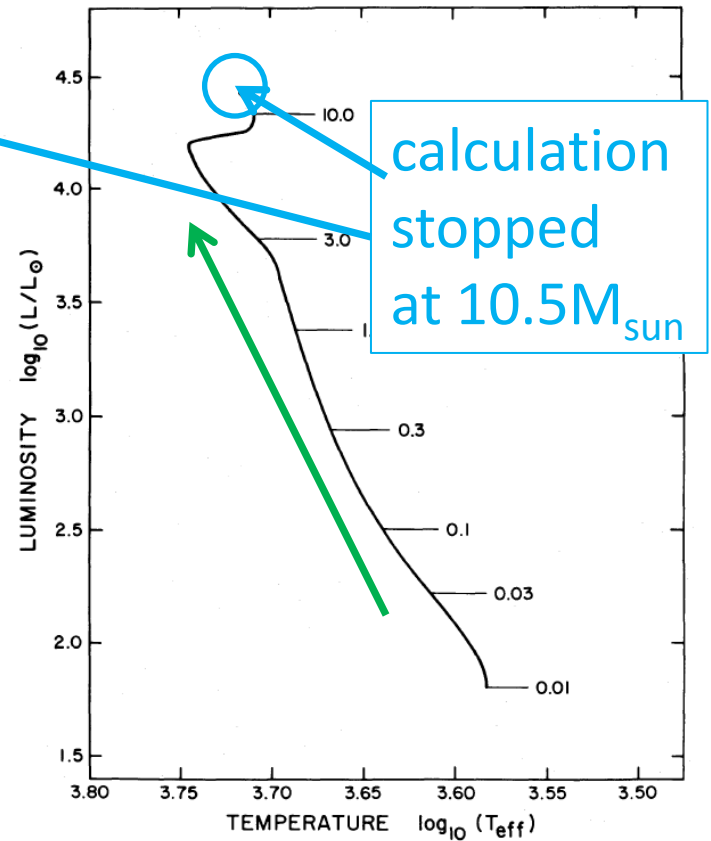
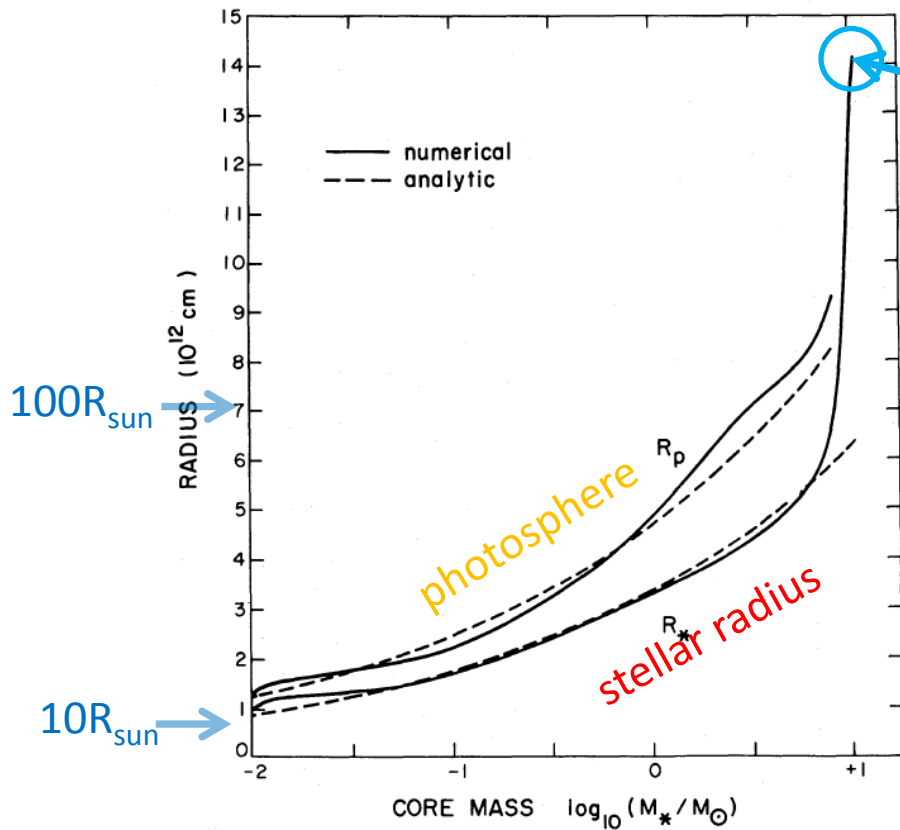
accretion rate chosen: $dM_*/dt = 4.4 \times 10^{-3} M_{\text{sun}}/\text{yr}$

protostellar radius:

much larger (several $10 R_{\text{sun}}$)
than present-day counterpart.

HR diagram:

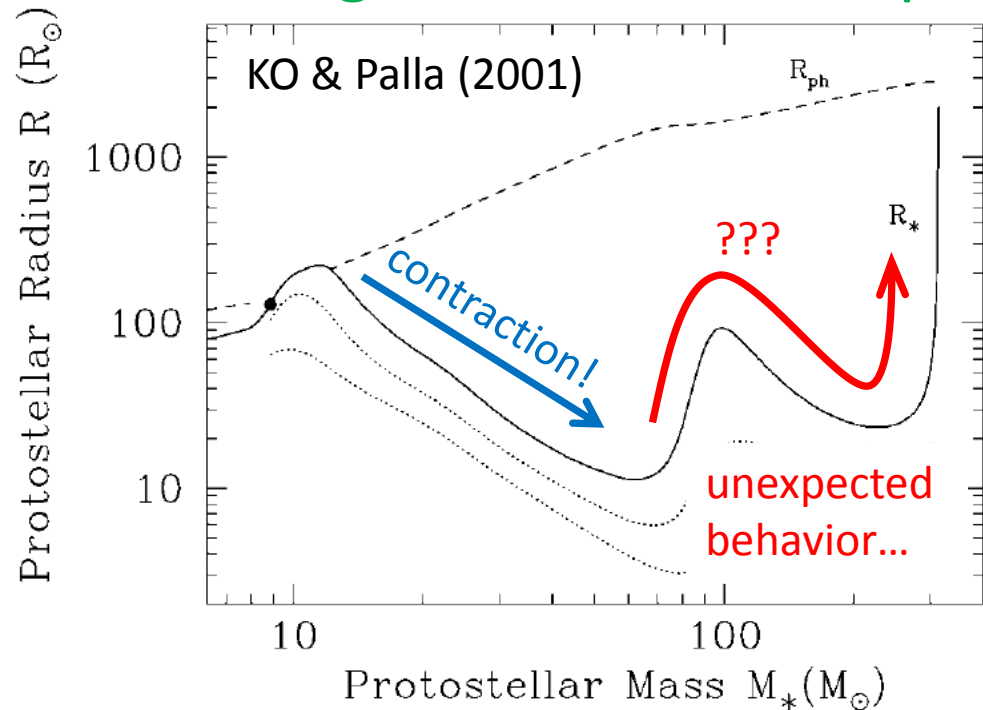
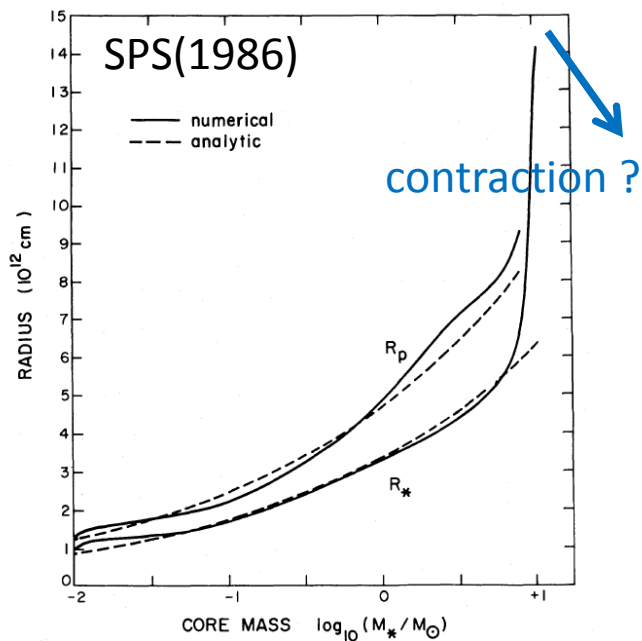
The protostar evolves almost
vertically with $T_{\text{eff}} \sim 5000\text{K}$



How is the evolution at $>10.5M_{\text{sun}}$?

Although we terminated the calculation at $M_* = 10.5 M_{\odot}$, it is clear from the trend of increasing internal luminosity that a fourth phase of rapid core contraction must ensue for higher masses. During this phase, the entropy of the deep interior will drop substantially while the central temperature rises. Eventually, hydrogen will be ignited in the central region.

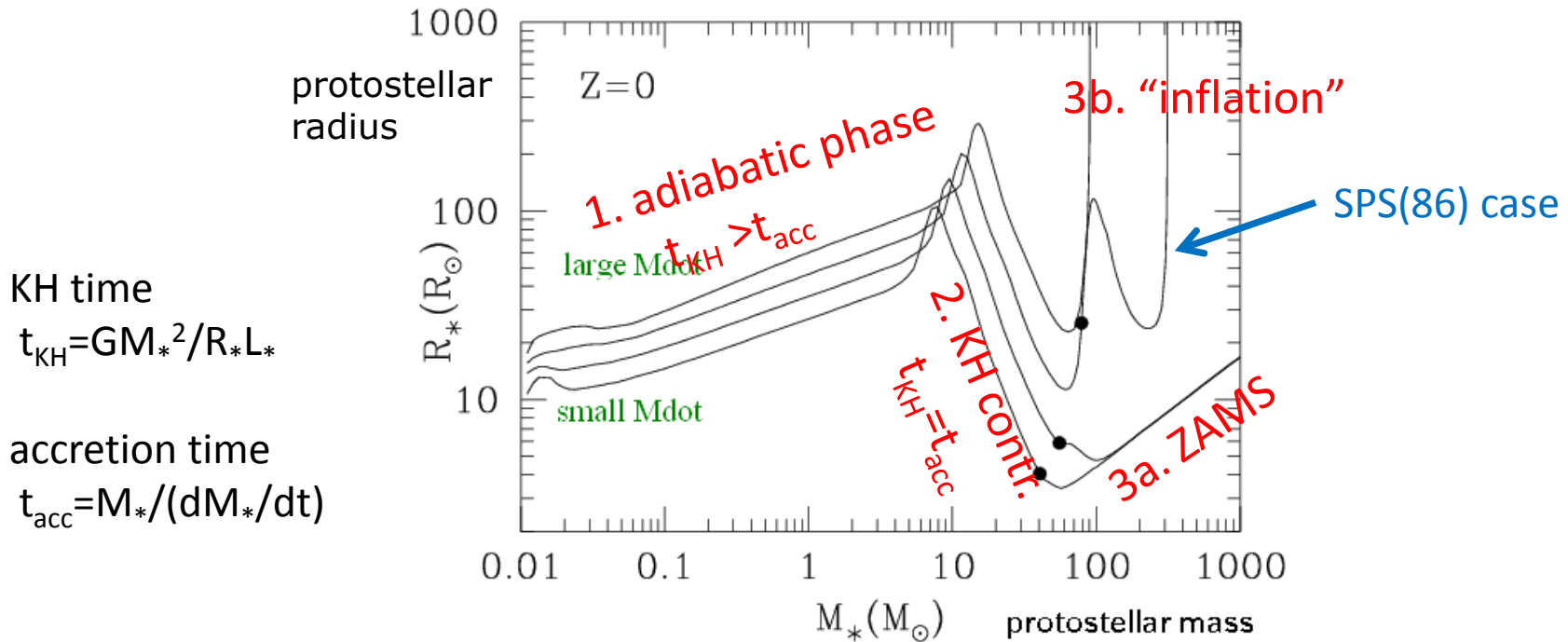
Yes, it was correct. But things were not so simple....



Cases with different accretion rates

KO & Palla (2003)

$$dM_*/dt = 8.8, 4.4, 2.2, 1.1 \times 10^{-3} M_{\text{sun}}/\text{yr}$$



All protostars go through the adiabatic and the KH contraction phases.

Subsequent phase depends on the accr. rate.

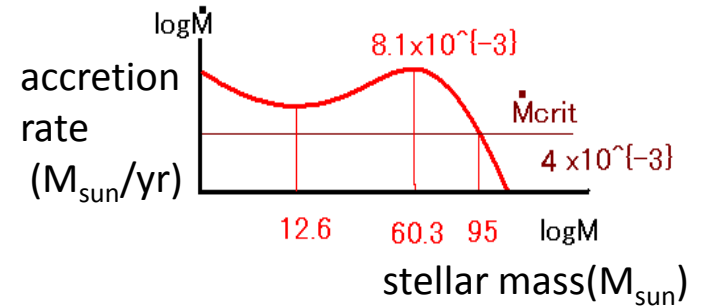
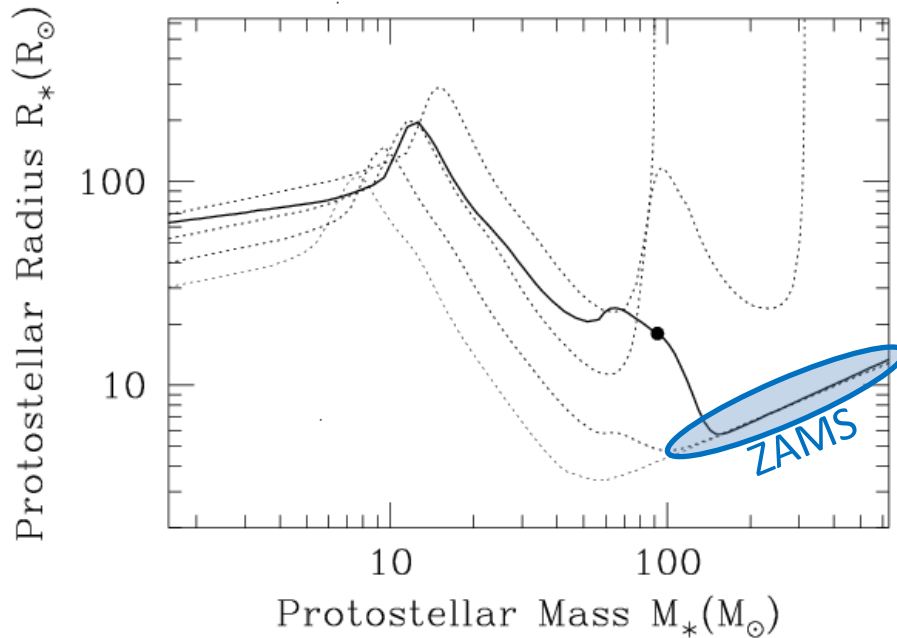
✓ With low accretion rate ($< dM/dt_{\text{crit}} = 4 \cdot 10^{-3} M_{\text{sun}}/\text{yr}$):
 → the star reaches ZAMS and accretion continues

✓ With higher accretion rate ($> dM/dt_{\text{crit}}$):
 → the star starts inflating when $L (= L_* + L_{\text{acc}})$ becomes close to L_{Edd} .
 (the end of accretion?)

$$L_{\text{acc}} \propto \dot{M}_*/R_*$$

With a realistic accretion history

Abel+(2002) accretion rate



- The protostar reaches ZAMS at $\sim 100 M_{\text{sun}}$.
- Accretion continues w/o stellar inflation.
- Since (MS lifetime $\sim 2 \times 10^6 \text{ yr}$) $>$ (core free-fall time $\sim 3 \times 10^5 \text{ yr}$), most of the core material can accrete within stellar lifetime.

Very massive star (\sim a few $100 M_{\text{sun}}$) is formed
in spherical accretion case

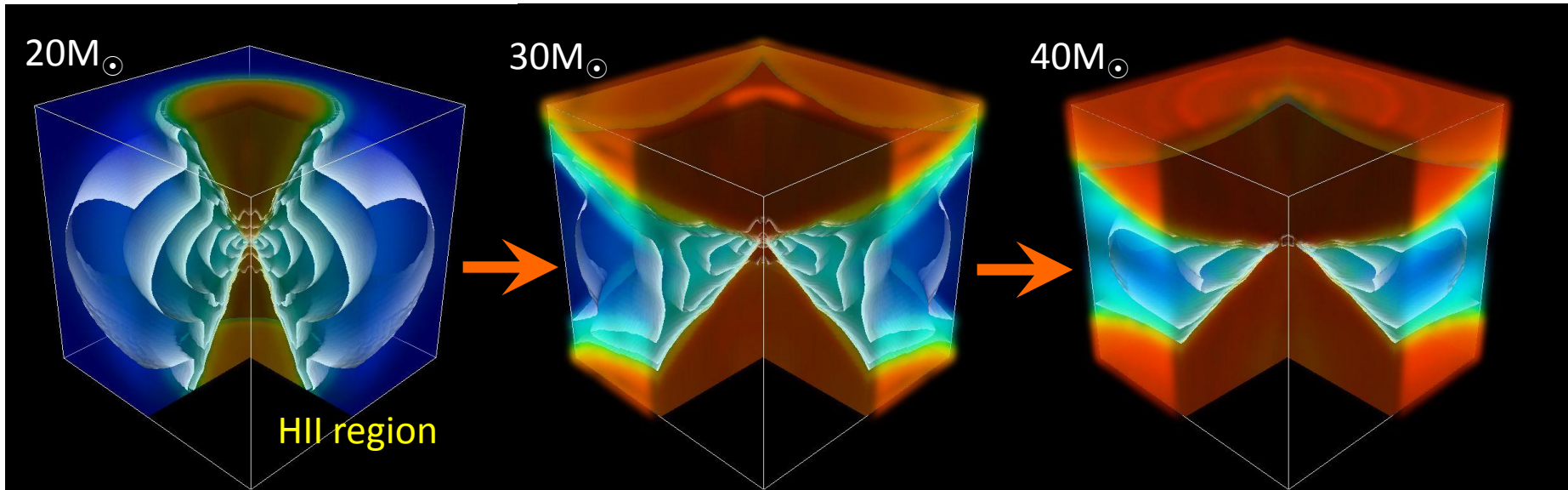
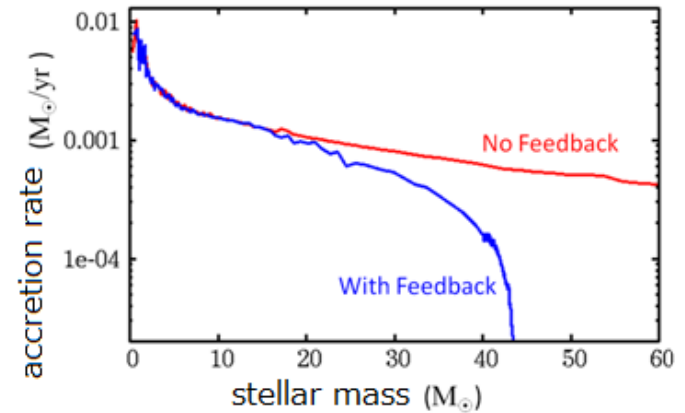
Feedback-limited mass of first stars

In non-spherical accretion:
mass of first stars is set by
the UV feedback

photoevaporation of disk

(McKee & Tan 08, Hosokawa+11/12,
Stacy+12, Hirano +14, Susa + 2014)

Accretion stops at $\sim 40M_{\odot}$



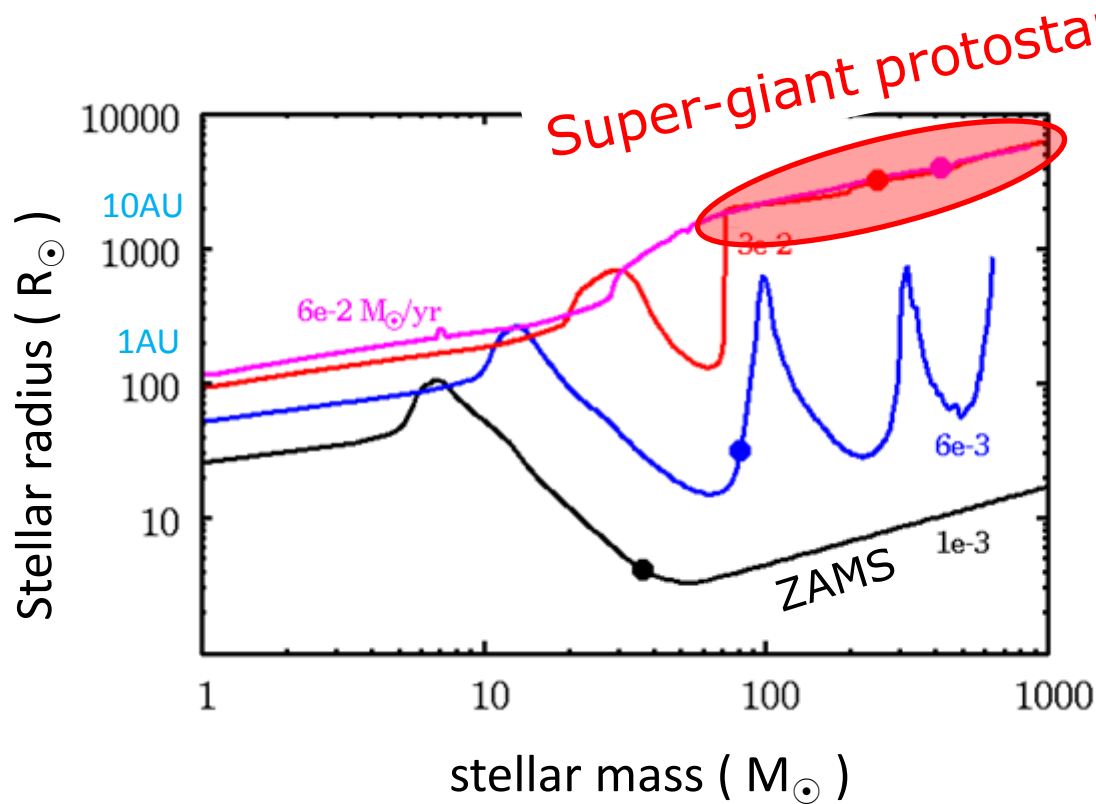
contour: density, color: temperature

Hosokawa, KO, Yoshida, Yorke 2011, 2012

See David Hollenbach and Takashi Hosokawa's talks

What if a very high accretion rate is maintained ?

What happens after protostar inflates?



- protostar starts inflating when $L \sim L_{\text{Edd}}$ with high accr. rate ($> 0.01 M_{\text{sun}}/\text{yr}$).

- stellar inflation stops at $\sim 10\text{AU}$ and star becomes super-giant without reaching main-sequence.

Hosokawa, Yorke, KO (2012)

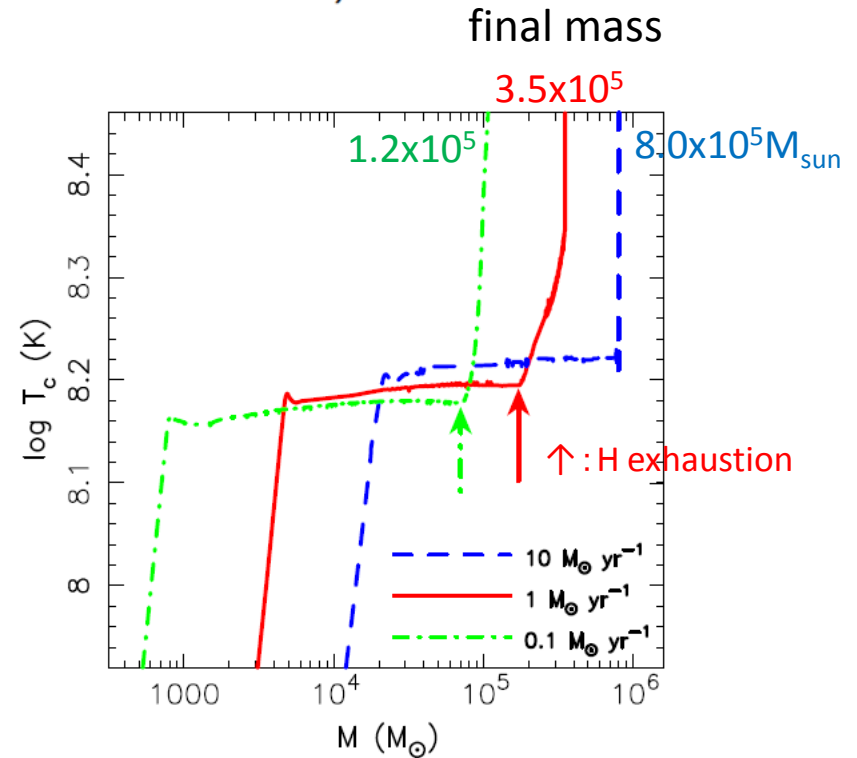
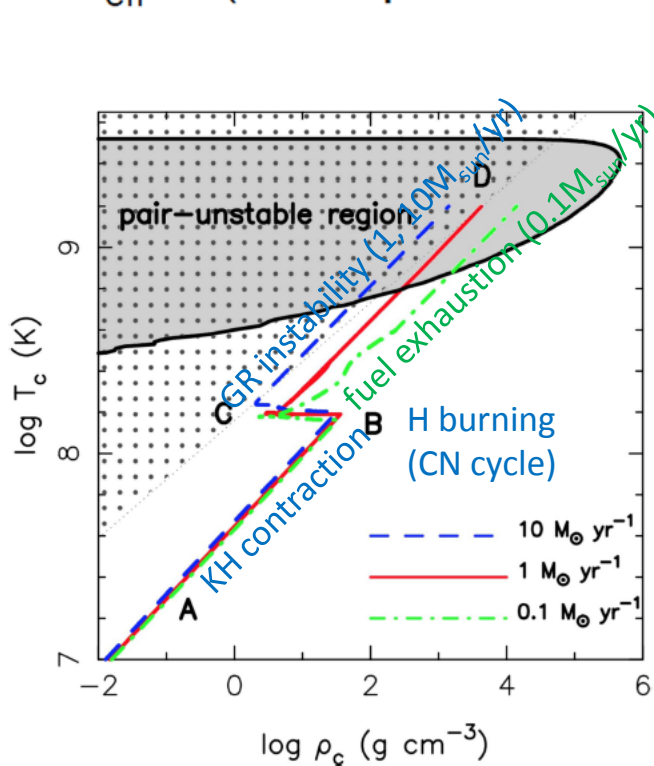
Supermassive star will be formed.
→ seed BHs for high-z SMBHs ?

General relativistic collapse

Accretion evolution with
Post Newtonian Gravity

Umeda, Hosokawa,
KO & Yoshida 2016

$$G_{\text{eff}} = G(1 + P/\rho c^2 + 4\pi r^3 P/Mc^2 + 2GM/rc^2)$$

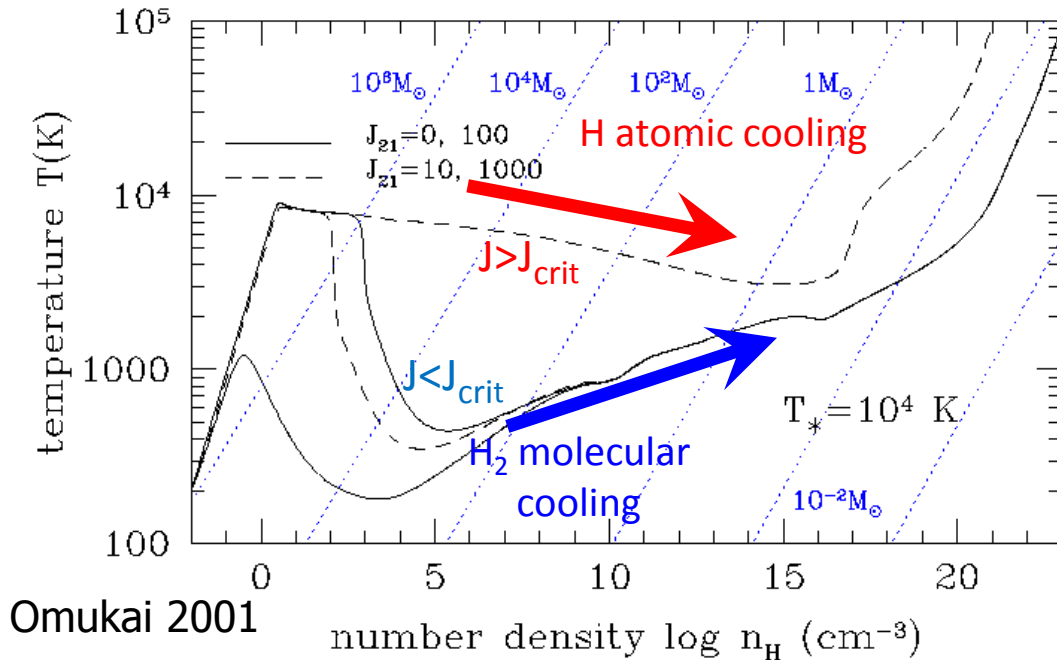


The star collapses at final mass of 10^5 - $10^6 M_{\text{sun}}$
depending on the mass accretion rate.

How such a high accretion rate is realized?

Possible pathway for super-massive stars: Atomic-cooling collapse

Bromm & Loeb (2003)

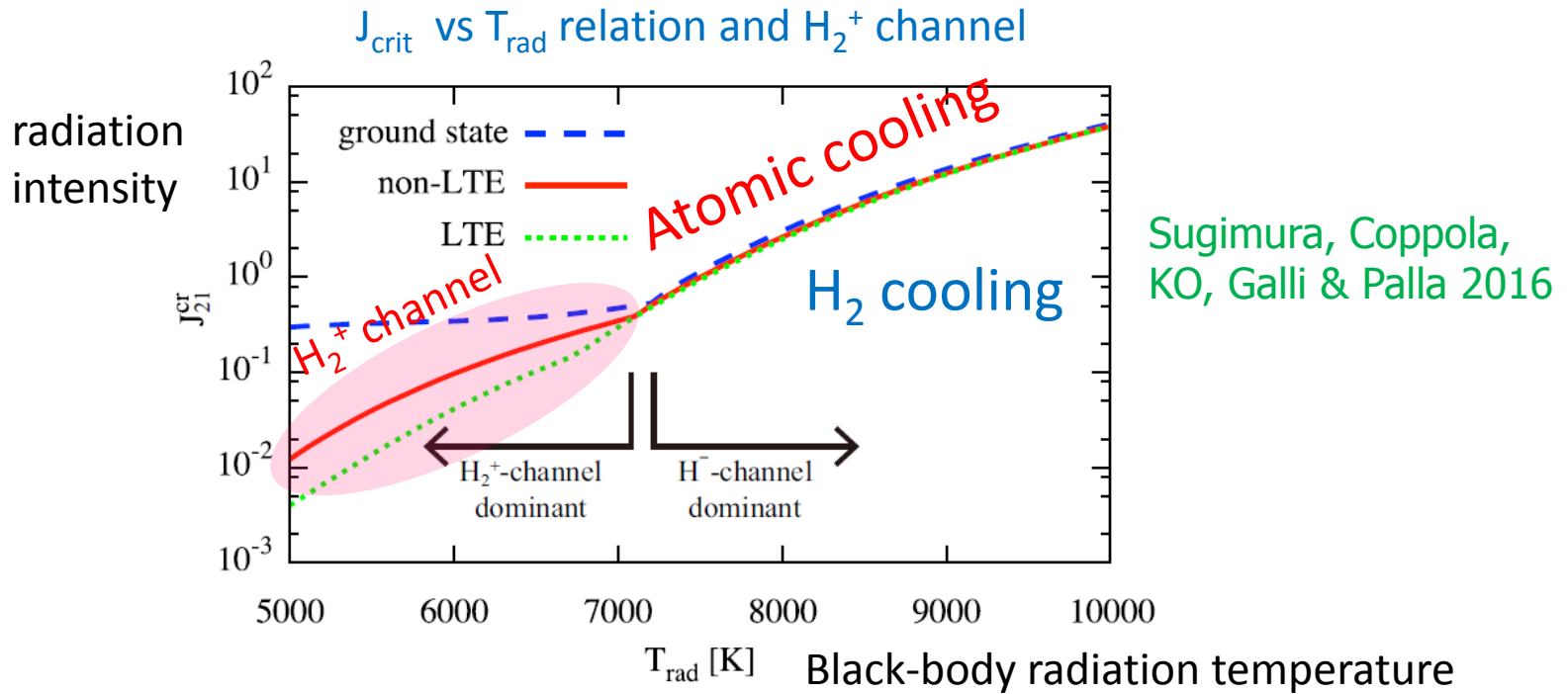


If FUV radiation is more intense than the critical value J_{crit} , the cloud cools solely by atomic cooling.

- high temperature (at $\sim 8000K$) during the collapse
→ high accretion rate in protostellar phase
 $dM_*/dt \sim 0.06 M_{sun}/yr (T/10^4 K)^{3/2}$
- no rapid cooling phase
→ monolithic collapse

See Dominik Schleicher's talk

How much FUV is needed for atomic-cooling collapse?



- H_2^+ channel (and so its non-LTE level population) is important for radiation field with very soft spectrum $T_{\text{rad}} < 7000\text{K}$.
- This justifies the previous estimate for J_{crit} for ordinary young galaxies, which have harder spectrum ($T_{\text{rad}} \sim 20000\text{K}$).

See Kazu Sugimura's poster (#104)

SUMMARY

The prestellar collapse of the first stars is rather well established:

- The gas becomes fully molecular by 3-body reactions at $\sim 10^{11}\text{cm}^{-3}$ and a small protostar is formed at the center at $\sim 10^{21}\text{cm}^{-3}$.

Multi-D simulations for the protostellar evolution are still underway:

- Protostar grows by a rapid accretion at $\sim 10^{-3}M_{\text{sun}}/\text{yr}$.
- Protostellar radiative feedback sets the final stellar mass.
- Mass distributions is probably biased to massive objects at a few 10- a few $100M_{\text{sun}}$

In some unusual circumstances, e.g. in strong FUV fields., supermassive stars might form by atomic cooling.

**Francesco has contributed greatly
in establishing all this picture**

Grazie di cuore Francesco!



First Stars IV, 2012, Kyoto